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Advanced
Photon
Source



Theory and Simulation of CSR Microbunching in Bunch Compressors

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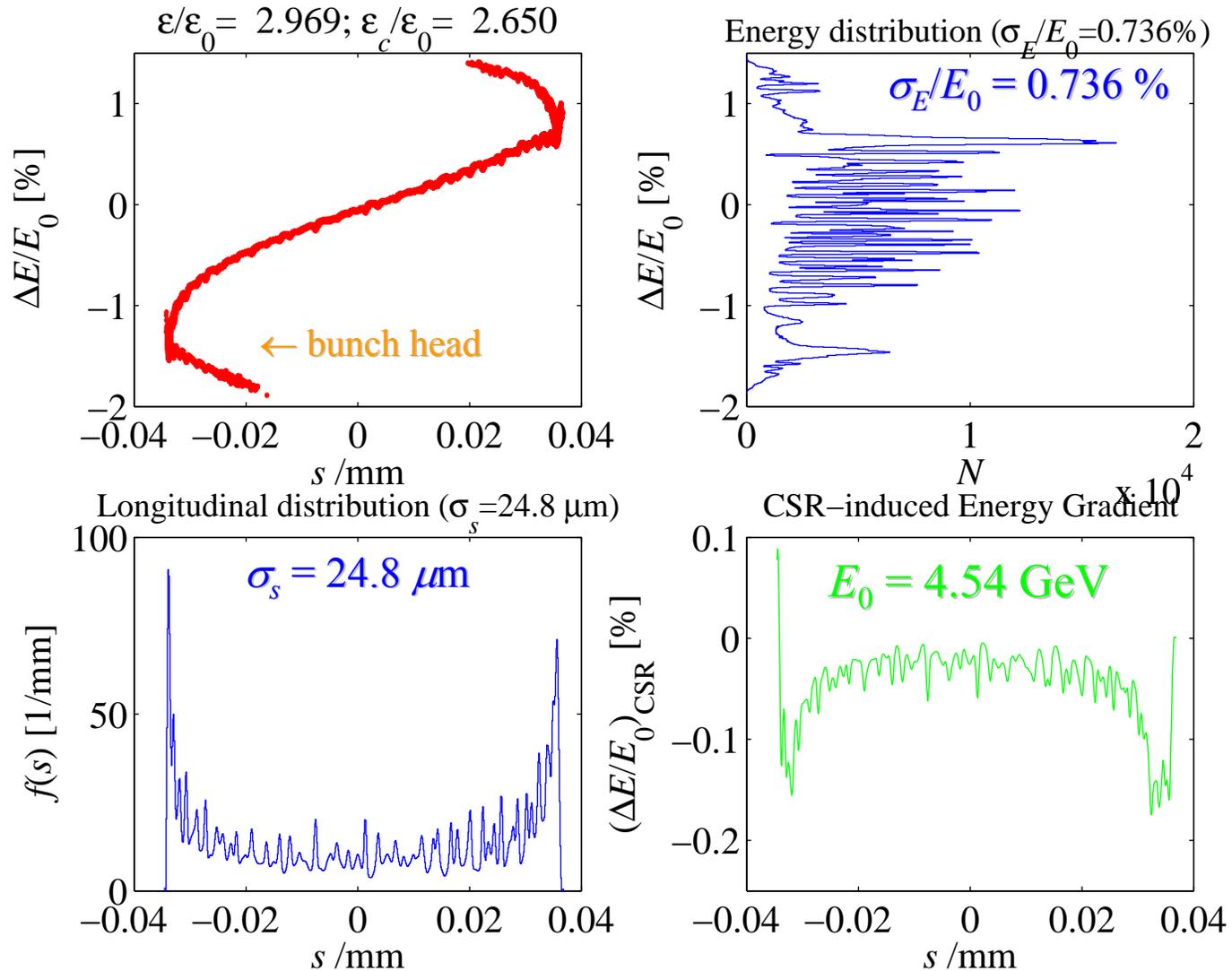
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Introduction

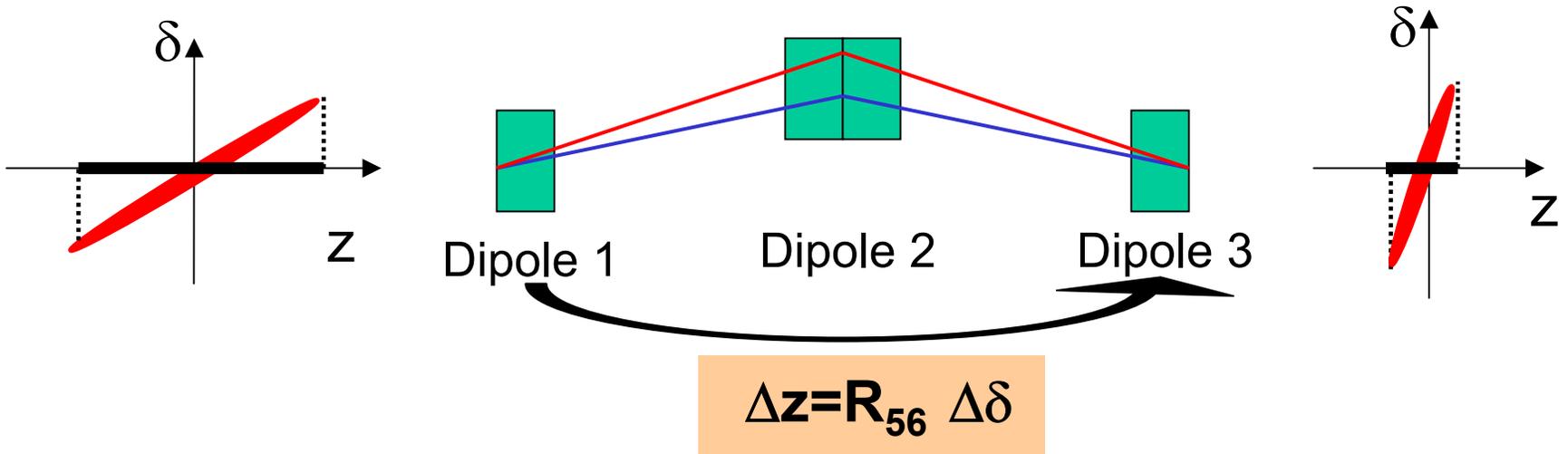
- X-ray FELs require very bright e-beams: low-emittance ($\sim 1 \mu\text{m}$), high peak current (\sim a few kA)
- PC-RF guns: low-emittance, low current ($\sim 100 \text{ A}$)
- Bunch length compression ($4 \text{ ps} \rightarrow 100 \text{ fs}$, $100 \text{ A} \rightarrow 4 \text{ kA}$), must not degrade the emittance
- CSR effects in compressor chicanes very challenging!
- Recently, a CSR microbunching instability is found by start-to-end LCLS simulations (Borland), could potentially corrupt the sliced emittance and energy spread

LCLS Distribution After BC2 Chicane

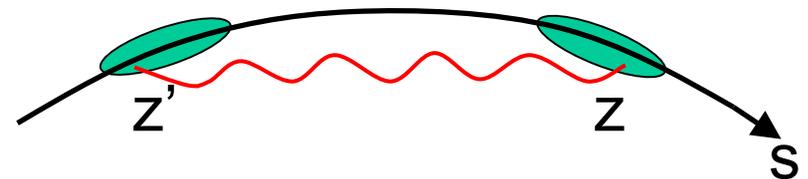


Compressor and CSR

chirped beam \rightarrow path length difference \rightarrow compressed beam



- Radiation from bunch tail catch up the head, increase energy spread and emittance

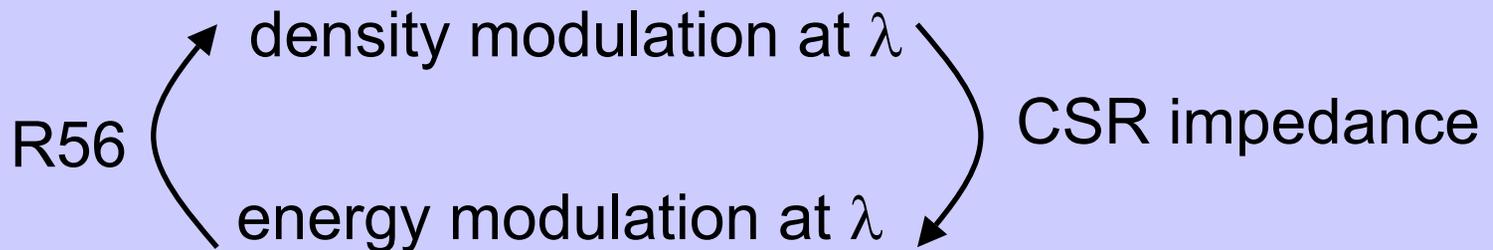


- CSR “wake” $W(z-z')$ or impedance $Z(k)$ ($k = 2\pi/\lambda$) (Derbenev et al., Murphy et al.)

CSR Microbunching

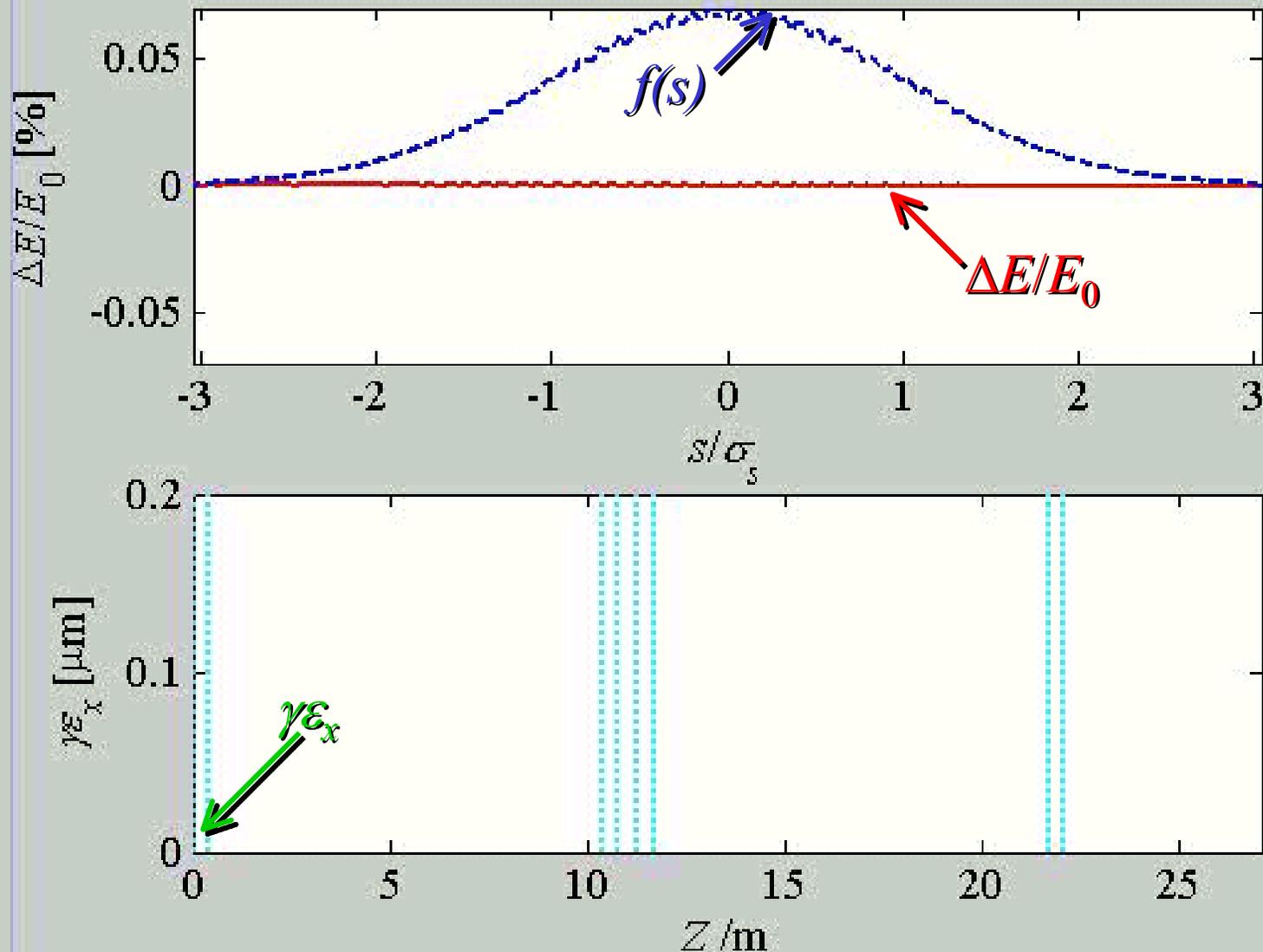
- CSR emitted from sub-bunch structures for wavelengths $\lambda \ll \sigma_z$ interacts back to the bunch, leading to a microbunching instability

(similar to microwave instability in a ring)



- This process can be initiated by either density or energy modulation

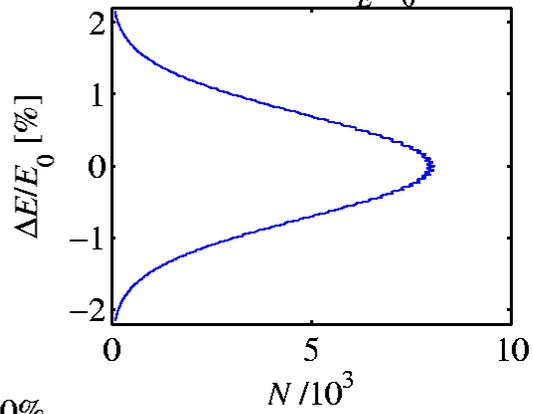
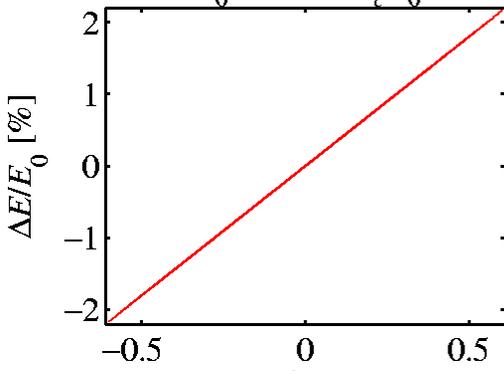
CSR Microbunching Simulation



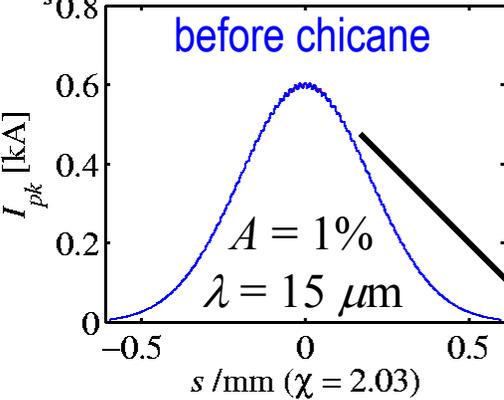
Initial: $\epsilon/\epsilon_0 = 1.000$, $\epsilon_c/\epsilon_0 = 1.000$

Energy distribution ($\sigma_E/E_0 = 0.7103\%$)

$\gamma\epsilon_x = 0$
 $\sigma_E/E_0 = 2 \times 10^{-6}$

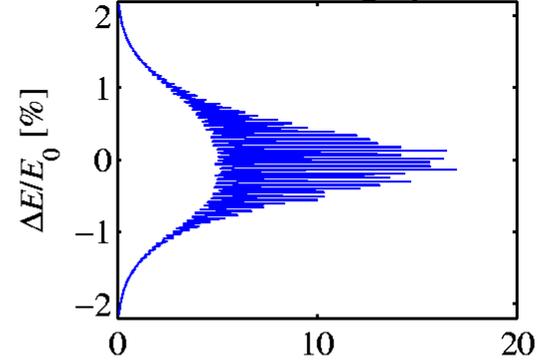
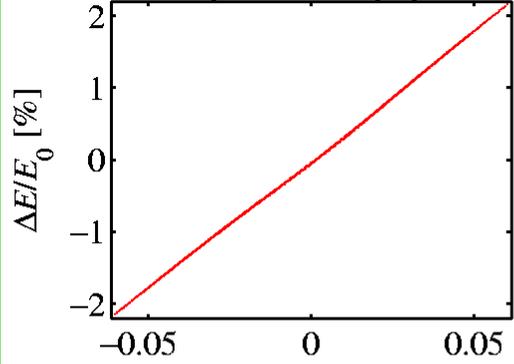


$\sigma_s = 0.1973$ mm, $\lambda = 14.998$ μ m, $A = 1.0\%$



Final: $\epsilon/\epsilon_0 = 30.465$, $\epsilon_c/\epsilon_0 = 25.659$

Energy distribution ($\sigma_E/E_0 = 0.7036\%$)

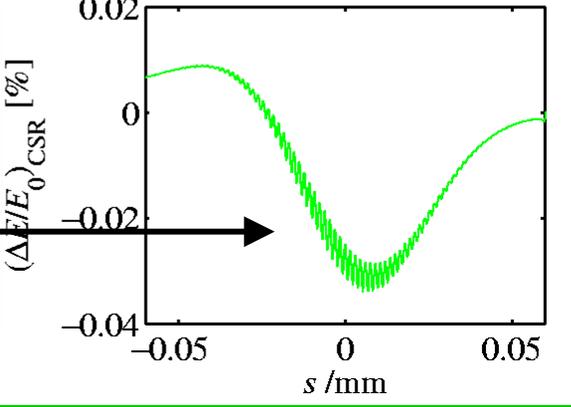
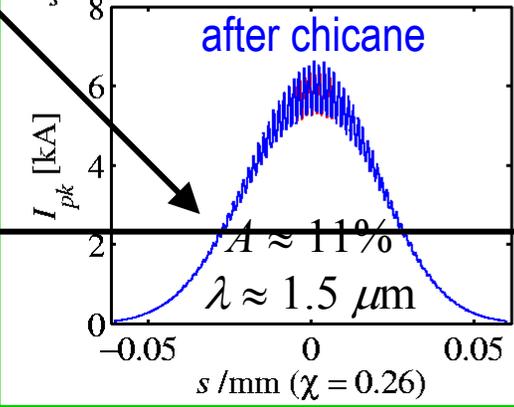


$\sigma_s = 0.0199$ mm, $\lambda = 1.525$ μ m, $A = 10.7\%$

CSR-induced Energy Gradient

$G \approx 11$

energy modulation



Theory

- Define a bunching parameter b at modulation wavelength $\lambda = 2\pi/k$ (fourier decomposes the current)
- Linear evolution of $b(k;s)$ can be described by an integral equation (Heifets et al.; Huang&Kim)

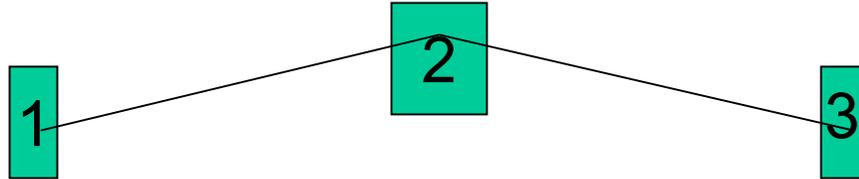
$$b(k(s);s) = b_0(k(s);s) + \int_0^s d\tau K(\tau,s)b(k(\tau);\tau)$$

$$\text{kernel } K(\tau,s) = ik(s)R_{56}(\tau \rightarrow s) \frac{I(\tau)}{\mathcal{I}_A} Z(k(\tau)) \times \underbrace{\exp(\dots\mathcal{E}, \sigma_\delta\dots)}_{\text{Landau damping}}$$

- For arbitrary initial condition (density and/or energy modulation), this determines the final microbunching

Staged Amplification

- Typical chicane



dipole separation $\Delta L \gg$ dipole length L_b

$$R_{56}(\tau \rightarrow s) \sim \begin{cases} \frac{L_b^3}{\rho^2} & \text{within the same dipole} \\ \frac{\Delta L L_b^2}{\rho^2} & \text{from one dipole to another} \end{cases}$$

- Ignore the induced bunching from energy modulation in the same dipole (Saldin et al.)
- Consider staged amplification from dipole to dipole by setting $K(s',s)=O(L_b/\Delta L)=0$ if $s-s' < \Delta L$

Iterative Solution

- Integral equation can be solved by two iterations

$$b(k; s) = b_0(k; s) + \underbrace{\int_0^s ds' K(s', s) b_0(k'; s')}_{\text{one - stage amplification}}$$

$$I_f(1 \rightarrow 3) + I_f(2 \rightarrow 3)$$

dominant in low-gain

$$+ \underbrace{\int_0^s ds' K(s', s) \int_0^{s'} ds'' K(s'', s') b_0(k''; s'')}_{\text{two - stage amplification}}$$

two - stage amplification

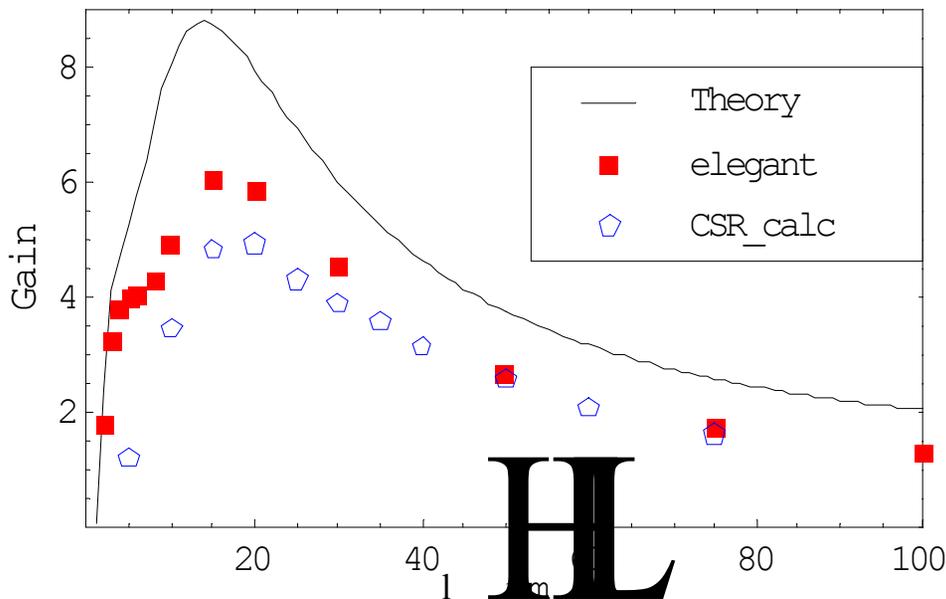
dominant in high-gain

$$I_f^2(1 \rightarrow 2 \rightarrow 3)$$

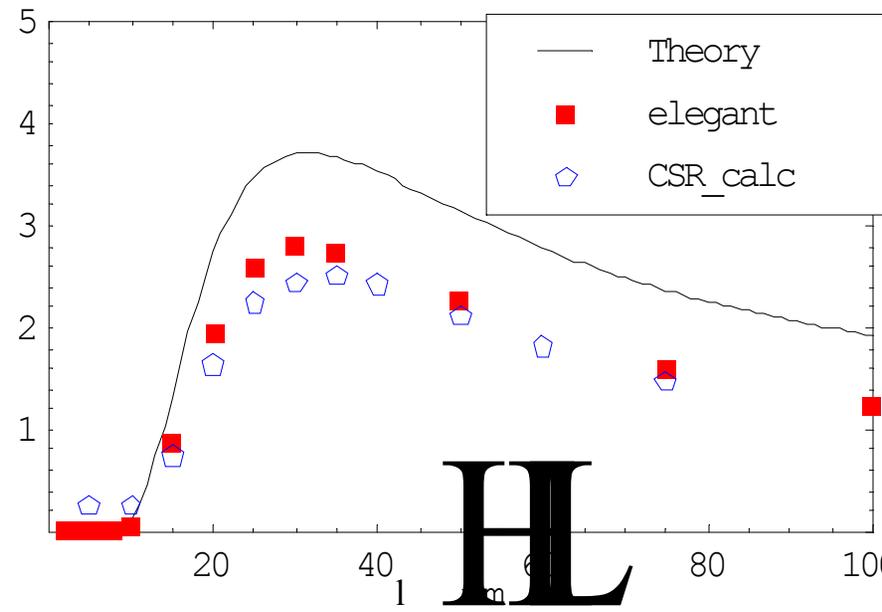
- Calculate gain= $|b_{\text{final}}/b_{\text{initial}}|$ as a function of λ , and compare with simulation results

Berlin CSR Benchmark Chicane

- Elegant and CSR_calc (matlab based) codes used
- a few million particles are loaded with 6D quiet start
- CSR algorithm based on analytical wake models

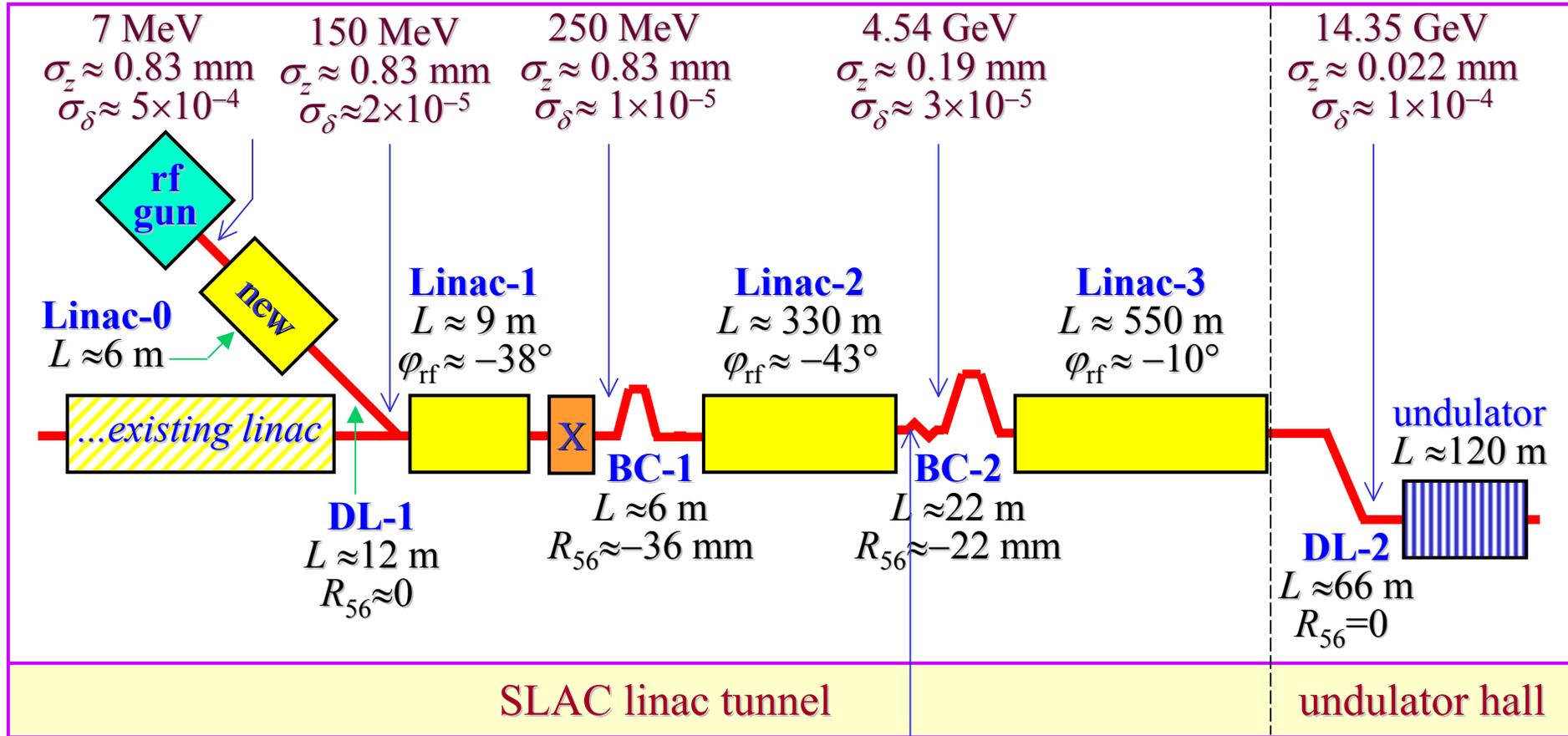


$$\sigma_\delta = 2 \times 10^{-6}, \gamma \epsilon_x = 1 \mu\text{m}$$



$$\sigma_\delta = 2 \times 10^{-5}, \gamma \epsilon_x = 1 \mu\text{m}$$

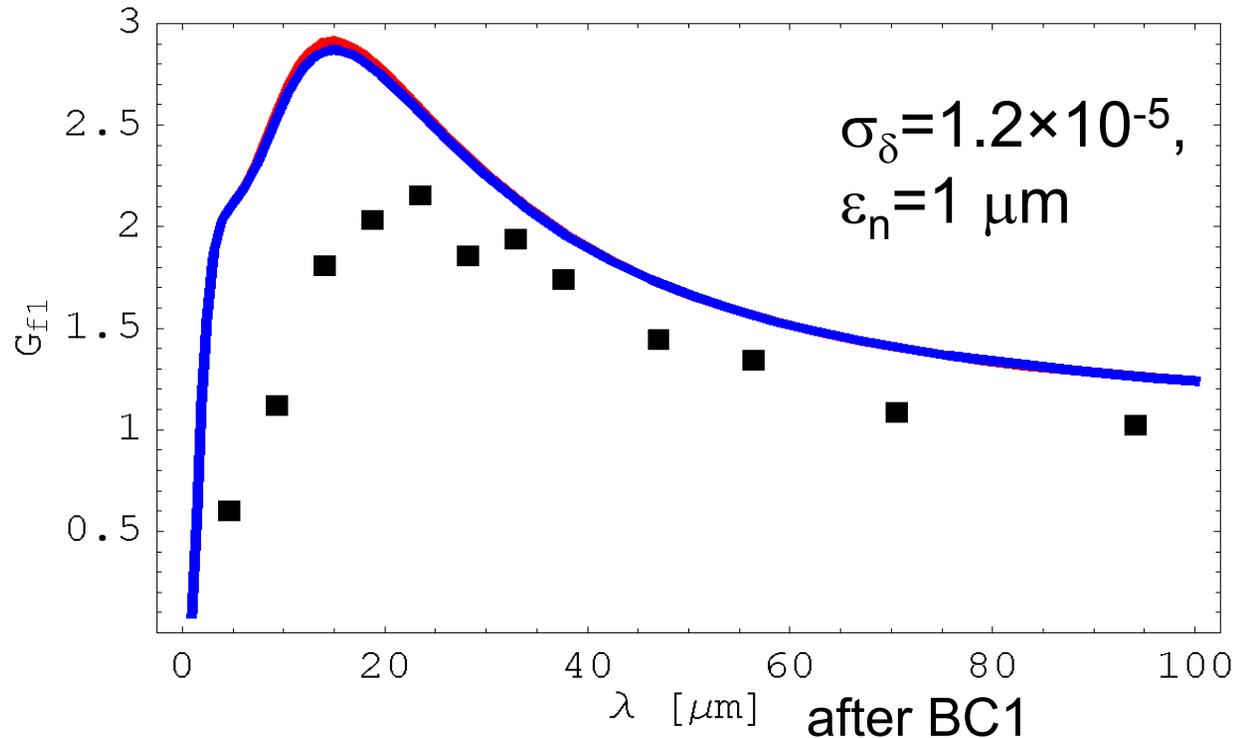
LCLS Acceleration and Compression Systems



To damp the instability, a SC wiggler can be placed right before BC2
 To increase the incoherent energy spread (still small for FEL instab.)

BC1

- BC1 gain in density modulation is low (so is BC2)

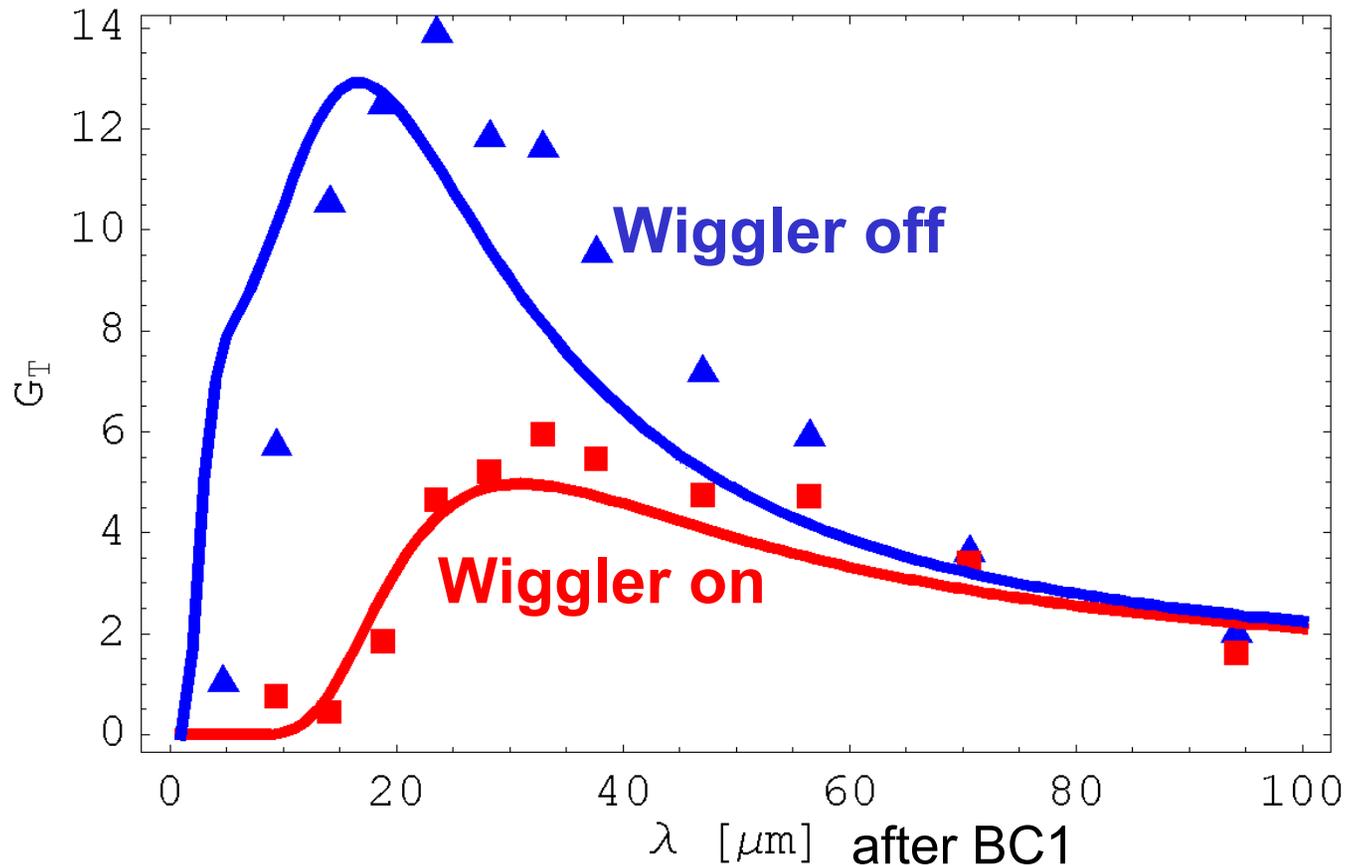


- CSR also induced energy modulation in BC1

$$\Delta p(s) \approx -\int_0^s ds' \frac{I(s')}{\mathcal{H}_A} Z(s') b(s') \exp(\dots \epsilon, \sigma_\delta \dots)$$

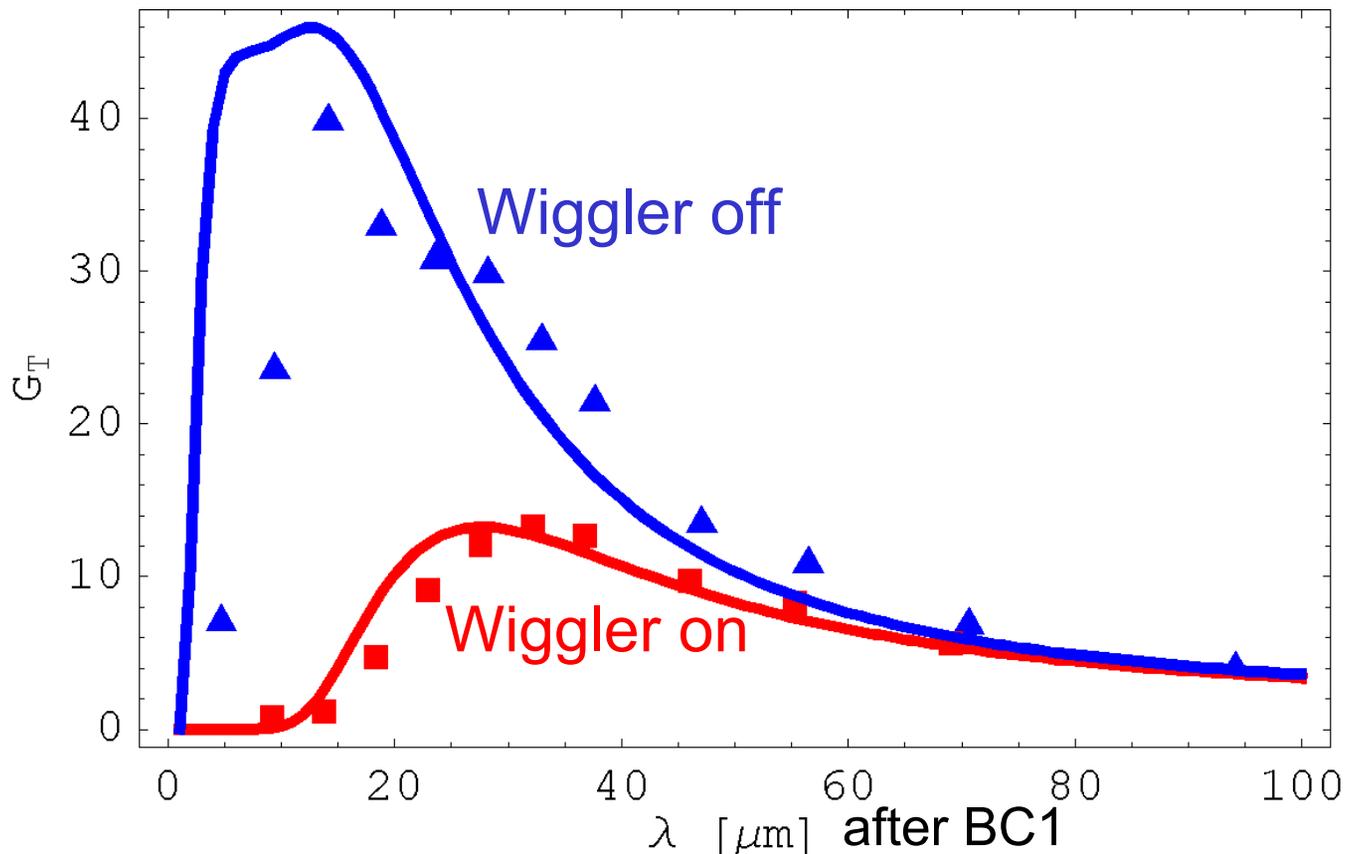
BC1+BC2

- BC2 not only amplifies density modulation gained in BC1, but turns BC1 energy modulation into gain in density modulation
- Total gain of BC1+BC2 > BC1 gain X BC2 gain



Full LCLS Gain

- LCLS has four bend systems: BC1, BC2, and two beam transport dog legs (DL1, DL2)
- DLs have very small R56 \rightarrow ignore induced density modulation but keep track of induced energy modulation



Comments on Initial Conditions

- From shot noise

$$b_{\text{eff}} \sim \frac{1}{\sqrt{N_{\text{coh. length}}}} \sim (10^{-4} \rightarrow 10^{-5})$$

with a gain less than 100, this should be a small effect

- From sharp bunch structure

(due to laser ripple in PC RF guns?)

$$b \approx \frac{N_{\text{spike}}}{N_{\text{total}}}$$



- Other sources of energy modulation (geometric wake...)
(simulations show more gain (2X) for full lattice and wake)
- Watch out for numerical noise in simulations!

Summary

- CSR microbunching instability is governed by an integral equation which is solved iteratively for chicanes
- Initialized by density and/or energy modulation, cascading through multiple chicanes and bends
- Two similar but independent CSR codes shows reasonable agreement with each other and with theory
- Full LCLS gain is significant but not detrimental for shot noise start-up
- The damping wiggler reduces the microbunching gain and the sensitivity to the initial bunch distribution